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TITLE

THE RF POWER SYSTEM FOR THE CHOPPER/BUNCHER SYSTEM ON THE NBS-LOS ALAMOS RTM

AUTHOR(S:

L(loyd) M. Young and D(avid) R. Keffeler

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# LOS Alamos National Laboratory Los Alamos, New Mexico 87545

L. M. Young and D. R. Keffeler, AT-1, MS H817 Los Alamos National Laboratory, Los Alamos, NM B7545 USA

#### <u>Summary</u>

The rf power system and its closed-loop feedback contract for the racetrack microtron (RTM) chopper/buncher system are described. Measurements made on the response of the feedback system to external perturbations will also be reported.

#### Introduction

The 100-keV injector for the National Bureau of Standards (NBS)-Los Alamos racetrack microtron (RTM)1 uses two square deflection cavities and a bunche: cavity. The deflection cavities operate in the  $TE_{102}$  and TE201 modes, and the buncher operates in the IM010 mode. Figure 1 illustrates how the beam is deflected off axis into a spiral by the deflection cavity; the Leam traces a circle on the chopping aperture. A 60° circumferential slit in the chopping aperture allows only 1/6 of the input beam to pass through. A pair of thin-gap lenses centered on this aperture focuses the chopped beam back to the axis at the second de flection cavity. The second half of the chopper sec tion is a mirror image (centered on the aperture) of the first half. If rf phase and amplitude are ad justed properly, the transverse momentum imparted to the beam by that of the first deflection cavity will be exactly cancelled by that of the second deflection cavity so that the beam becomes coaxial after the second deflection cavity.

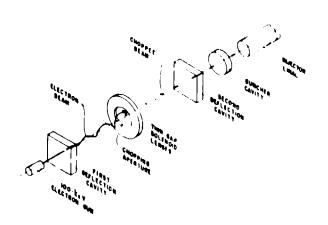


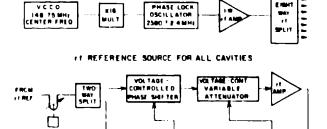
Fig. 1. The chopper/buncher system for the HS Los Alamos RIM An elactron beam from a 100-keV electron gun is deflected by a rotating magnetic field in the first deflection cavity, forming a spiral that traces a circle on the chopping aperture. The thin-gap lenses focus the spiraling beam back to the axis at the second deflection cavity. There, it is deflected along the axis by a rotating magnetic field that is adjusted to exactly cancel the transverse kick it was given by the first deflector.

The 2380-MHz bunching cavity decelerates the head of the 60° long bunch from the chopping system and accelerates the tail. This action reduces the length of the bunch to about 10° at the entrance to the injector linac. The square deflection cavity deflects the beam into circles on the chopping aperture by simultaneous excitation of the 11.102 and 11.201 modes. One mode deflects the beam horizontally and the other vertically. When the modes are excited to equal amplitude with relative rf phase difference of 90°, the beam scribes a circle on the chopping aperture. This chopping system causes very little emittance growth of the

electron beam, but only if the phases and amplitudes of the rf power in each deflection cavity is controlled very accurately. Therefore, each deflection cavity uses two rf power sources, each with its own phase and amplitude control. The chopper/buncher system uses a total of five rf power sources with five separate amplitude and phase feedback-control loops. These amplitude and phase feedback loops are all identical.

#### Phase and Amplitude Feedback Controls

The rf system is shown in Fig. 2. The rf source is derived from a voltage-controlled crystal oscillator (VCCO) at 148.75 MHz, which is then frequency multiplied by 16 to 2380 MHz. The frequency of this VCCO can be adjusted ±0.1% with a ±5-V control signal. The rf source has phase noise from the multiplication; therefore, an oscillator is phase locked to this source. The output of this phase-lock oscillator has very little phase noise.



TYPICAL 11 FEEDBACK-CONTROL SYS IEM FOR RTM

Fig. 2. A common of reference source is used for all the cavities in the RIM. Each of the eight outputs of the moreference source is used as an input to a feedback control system. Tack cavity has its own feedback system.

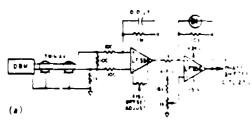
A good rf source without phase noise is required for the RTM because there are many cavities in the RTM and all must operate at constant phase with respect to each other. A noisy of reference source would prevent accurate phase control and measurements. An eight-way splitter on the rf source provides the rf reference for all the rf control loops in the RIM. These rf reference signals go through stepping-motor-driven line stretchers to provide phase adjustment for each device. In each controller, the rf reference is split with one output going to a double balanced mixer (DBM) and the other going through a voltage-controlled phase shifter and a voltage controlled attenuator to one or more amplifiers required to drive a particular cavity. The voltage controlled phase shifter has a phase range of -270° for a control voltage swing from 0 to 30 V. The voltage controlled attenuator has a range of ~60 dB, with an attenuation change of ~10 dB/V from 0 to 6 V. These devices are not perfect in that the phase shifter has a small change of attenuation versus control voltage, and the attenuator has a small amount of phase shift versus control voitage. The small in teraction between amplitude and phase control has not caused any problems with operation and control of the chopper buncher.

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A pickup loop in the cavity samples the rf cavity fields close to the outer wall. This rf signal is split with one output going to an amplitude detector (low-barrier Schottky diode detector). The output of this detector is proportional to the rf amplitude in the cavity. The other output goes to the DBM, where it is mixed with the rf reference. The output of the DBM will consist of two frequencies: the sum and the difference. The sum frequency can easily be removed with a low-pass filter. Because the two inputs to the DBM are the same frequency, the frequency of the difference is zero, which is a dc signal. This dc signal provides the phase information and can be represented by A x sin  $(\Delta \varphi \cdot \varphi_{C})$  , where A is a function of the rf power of the two signals, ap is the phase difference of the two signals, and  $\phi_C$  is a constant. By choosing to operate the phase-control loop with the output of the DBM zeroed,  $\Delta \phi$  must be equal to  $\phi_C$  + n x 180°, where n is an integer and  $A \neq 0$ . The voltage signals from the amplitude detector and the phase detector (the DBM) are amplified by the feedback amplifiers. The amplified signals are then used to drive the phase shifter and the variable attenuator.

The operational amplifier (DP-AMP) chosen for the feedback amplifiers is the LF356N. It has a gain bandwidth of 4 MHz and a low-input offset voltage temperature coefficient of  $\leq 3~\mu\text{V/°C}$ , which results in a low-input offset drift that is important for good stability of the amplitude and phase. Figure 3 shows a simplified schematic of the feedback amplifiers. All the OP-AMPs are in an inverting configuration. The phase feedback amplifier requires only two OP-AMPs. The first is the integrating stage, and the second provides a voltage output between 0 and 23 V. The O-to 23-V range utilizes most of the 270° of phase shift available from the phase shifters because the phase shifters are near saturation for control voltages be tween 23 and 30 V.

Because the output of the DBM is bipolar (with two zero-crossing points) versus phase difference be tween the rf reference and the rf sample, one of the crossing points will result in positive and the other in negative feedback. The circuit automatically will



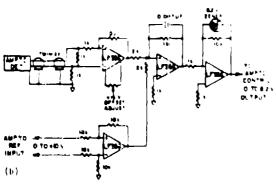


Fig. 3. The feedback amplifiers for (a) the phase control and (b) the amplitude control use the same methods to reduce noise pickup on the input. The DBM and the implitude detector are only grounded by the shield of the twin as of the analog ground of the feedback amplifiers. The DBM and the amplitude detectors are iso lated from the ground of the rf system by dc blocks (not shown)

select the appropriate zero-crossing point that results in negative feedback for the phase control. The zero crossing that results in positive feedback is unstable, and the circuit will try to change the phase to the nearest zero crossing with negative feedback. It is possible to set up the phase-control loop with the two zero-crossing points of the DBM within the 270° range of the phase shifters. Then, it is possible for the phase-control loop to latch on the wrong side of the zero-crossing point with positive feedback. To eliminate this possibility, a stepping-motor controlled phase shifter is included in the phase-control loop to adjust the phase-control range to encompass only the negative-feedback, zero-crossing point.

The cutput of the amplitude detector is unipolar; therefore, the polarity of the feedback is fixed. The amplitude detector has a negative output, and the variable attenuator increases the attenuation for a positive control voltage, which requires the feedback amplifier to be inverting. In this system, the best way to achieve the inversion is with three inverting OP-AMPs, consisting of an input buffer, an integrator, and an output buffer.

An example of a simple feedback-control system is shown schematically in Fig. 4(a). Figure 4(b) shows this system open loop and Fig. 4(c) shows the Bode approximation of the open-loop, gain-magnitude frequency response. Figure 4(d) shows the corresponding Bode approximation of the phase shift. The open-loop gain must be less than 1 when the phase shift exceeds  $180^{\circ}$  or the feedback loop [Fig. 4(a)] will oscillate. In this example, A is an operational amplifier with one pole at frequency  $F_{\rm D}$ . This pole introduces a phase shift that equals  $90^{\circ}$  for all frequencies greater than  $\sim 10~F_{\rm D}$ . The load will introduce additional phase shift that causes the total phase shift to exceed  $180^{\circ}$  at a high frequency.

The loop gain of the system must be less than I when this additional phase shift causes the total phase shift to exceed 180°. The phase margin is defined as the amount the phase shift is less than 180° at the frequency when the loop gain equals 1.

The feedback amplifiers are configured as integrators only. Because the gain response must fall off uniformly at 20 dB/decade (resulting in a 90° phase margin in the feedback amplifier) and must have a unity-gain crossover that provides a satisfactory phase margin, the integrator capacitance is chosen to

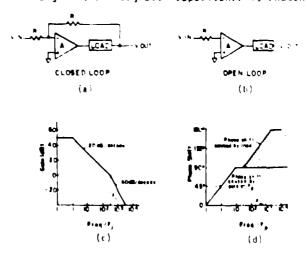


Fig. 4. The basic components of a feedback system are shown in the normal closed loop configuration, (a). If the feedback loop is open (b), the Bode approximation of the gain versus frequency is shown in (c), and the phase shift of the system is shown in (d). The amplifier A has a pole at frequency ip and the load has a pole at frequency if A feedback system is stable if the total phase shift in the feedback loop is less than the for frequencies at which the gain is greater than unity. If this condition is violated, the system will oscillate.

give a unity-gain crossover at the desired frequency. The feedback amplifiers are very easy to adjust because only one variable needs adjusting and that one is the frequency at which the loop gain is unity. The optimum size of the integrator capacitor depends on the gain of the loop external to the feedback amplifier. Because the variable attenuator is linearized, the feedback gain in the amplitude control loop is proportional to the magnitudes of the amplitude detector signal. The feedback gain in the phase-control loop also depends on the amplitude of the rf and the magnitude of the phase-control voltage. The phase shifter is not linearized and has more gain for phasecontrol voltages near O V than for control voltages near 30 V. The integrator capacitor used in the feed back amplifier gives a unity-gain crossover at ~40 kHz (see Fig. 5) for both control loops when there is a 0.4-V signal from the amplitude detector, which corresponds to an amplitude reference voltage of 2 V.

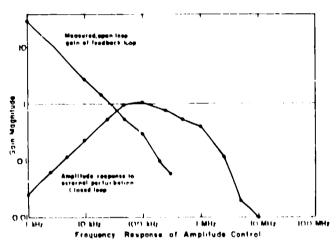


Fig. 5. Measurements of gain versus frequency with the amplitude feedback system operating open loop are indicated. With the feedback system operating in a closed loop configuration, a perturbation was applied to the load. The resulting ratio of the change in the rf amplitude to the amplitude of the perturbation versus frequency is shown.

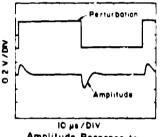
Figure 5 shows the open loop gain of the ampli tude feedback loop, which was measured with an ~0.3 V signal level from the amplitude detector. Note that the unity gain crossover occurs at ~30 kHz. For a 0.4 V signal, the unity gain crossover occurs at ~40 kHz, and so on. The amplitude response was meas ured with the loop closed and with a 0.4 V signal from the amplitude detector. This measurement was made by inserting an external, voltage controlled attenuator in the rf feedback loop and comparing the response of the rf amplitude signal with the feedback loop in control to the response without the edback loop in control. If the application required better control for perturbations near 100 kHz, higher speed OP-AMPs would have to be used in the feedback amplifiers; however, for the RIM this should not cause any difficulty be cause there should not be any source of perturbation at this frequency. Good control is obtained at all other frequencies. The open loop gain of only the feedback amplifier was measured and compared to the open loop gain calculated with the Bode approximation. The gain drops off slightly faster than the Bode an proximation predicts from the OP AMP specifications. this measurement of the open-loop gain indicates that in this circuit, the OP AMPs have a unity gain band width of ~3 MHz.

The measured open loop gain of the phase control loop essentially is identical to the gain of the amplitude feedback loop shown in fig. 5. This measurement was made with a 0.3 V signal on the amplitude detector and with the phase control voltage near the middle of

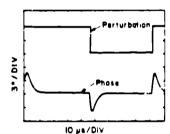
the control range. The gain is higher for the control voltage near zero and lower for the control voltage near 23 V. The phase response was measured by inserting a voltage-controlled phase shifter in the rf feedback loop and modulating the voltage on this external phase shifter. The response measurement was then made by comparing the phase perturbation with and without the phase feedback control working. The phase response is also essentially identical to the amplitude response shown in Fig. 5. This measurement shows that at 100 kHz, the feedback control has little effect on the perturbation, but that good control is obtained at all other frequencies. The response drops off above 100 kHz for both amplitude and phase because the high-Q cavity attenuates the perturbation.

Figures 6 and 7 show the amplitude and phase response to a square-wave perturbation. The perturbations are shown on the top trace, and the resulting effect on the amplitude and phase are shown on the bottom traces of Figs. 6 and 7, respectively. The scale on the perturbation and the response are the same. The perturbation causes an unwanted change in the amplitude and phase, but the amplitude and phase excursion is only about half that of the perturbation, and the amplitude and phase are returned to their desired value within ~10 us by the feedback controls.

Preliminary tests show that this rf system controlled the amplitude to within 0.1% and the phase to within 1°. Short-term variations in the amplitude and phase were <0.04% in amplitude and <0.1° in phase in a test <ystem.



Amplitude Response to Square-Wave Perturbation



Phase Response to Square-Wave Perturbation

Fig. 6. An oscilloscope trace of a square-wave perturbation and the resulting effects on the amplitude of the rf in a cavity. The size of the square-wave perturbation shown is adjusted to equal the resulting perturbation of the amplitude signal with the feedback turned off.

Fig. 7. An oscilloscope trace of a square-wave perturbation and the resulting effects on the phase of the rf in a cavity. The size of the square-wave perturbation shown is adjusted to equal the resulting change in phase with the feedback turned off

### Conclusion

This rf system was installed on the R1M at NBS in March 1984 and has operated satisfactorily during tests with beam in the chopper/buncher system. The electron beam was observed to scribe a circle on a view screen at the chopping aperture. The thin-gap lenses focused the beam back to the axis at the second deflection cavity. The rf fields in the second deflection cavity were adjusted to exactly cancel the transverse deflection it was given by the first deflection cavity. This cancellation of the transverse deflection was observed by a view screen beyond the second deflection cavity.

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